

# Selective Vibrational Mode Excitation in Nanocarbons by Tailored Intense Near-Infrared Pulses and Subsequent Intramolecular Vibrational Energy Redistributions

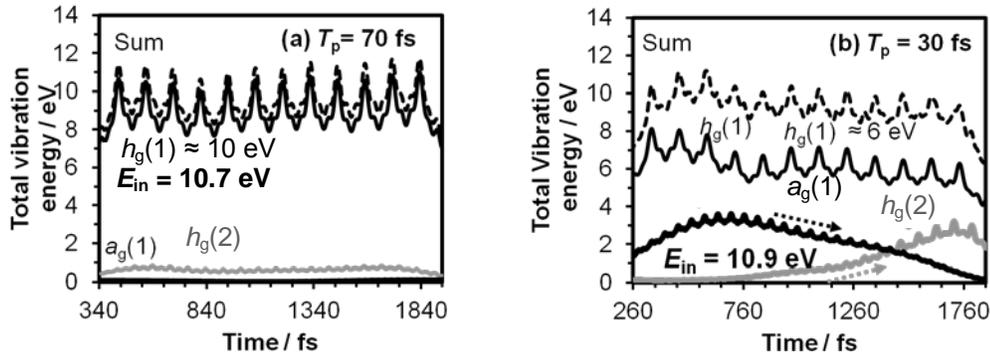
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Interactions between nanocarbons and intense near-infrared (NIR) pulses with femtosecond pulse length  $T_p$  and intensity  $I \approx 10^{15}$  W/cm<sup>-2</sup> induce structural rearrangements and fragmentations. Larmann *et al.* experimentally showed that tailored NIR pulse trains can optimize the specific fragment yields of fullerene C<sub>60</sub> [1]. We theoretically showed that selective impulsive Raman excitation of specific vibrational modes with vibrational period  $T_{\text{vib}}$  in C<sub>60</sub> can be achieved by setting  $T_p \approx T_{\text{vib}}/2$  and the fragmentation process depends on the vibrational modes initially excited [2]. These results suggest that intramolecular vibrational energy redistributions (IVR) leading to fragmentations depend on the pulse shape.

We performed the *on-the-fly* trajectory calculations on NIR-pulse-induced vibrational dynamics in C<sub>60</sub> with various  $T_p$  to clarify the effects of the pulse shape on the IVR dynamics. We combined the time-dependent adiabatic state approach [3] and density-functional based tight-binding theory [4] for simulations. We found that exciting different vibrational modes drastically changes the IVR dynamics as in Figs. 1a and 1b even if the total vibrational energy  $E_{\text{in}}$  acquired remains the same. Exciting the  $h_g(1)$  mode ( $T_{\text{vib}} = 140$  fs) by setting  $T_p = 70$  fs leads to the slow IVR in Fig. 1a; fast and complex IVR between the  $a_g(1)$  and  $h_g(2)$  modes occurs as in Fig. 1b when the  $a_g(1)$  mode ( $T_{\text{vib}} = 60$  fs) is enhanced by setting  $T_p = 30$  fs. We will also discuss IVR dynamics in a semiconducting single-wall carbon nanotube [5].



**Fig. 1:** Temporal change of the vibrational energies in low frequency modes of C<sub>60</sub> induced by single pulse excitation ( $I = 7.0 \times 10^{14}$  W/cm<sup>2</sup>,  $\lambda = 1800$  nm). (a)  $T_p = 70$  fs: the  $h_g(1)$  mode ( $T_{\text{vib}} = 140$  fs) is mainly excited and the IVR is slow. (b)  $T_p = 30$  fs: the  $a_g(1)$  mode ( $T_{\text{vib}} = 60$  fs) is relatively enhanced and fast IVR between the  $a_g(1)$  and  $h_g(2)$  modes is observed.

**References:** [1] T. Laarmann *et al.*, *Phys. Rev. Lett.*, 98:058302, 2007; [2] N. Niitsu, K. Yamazaki *et al.*, *J. Chem. Phys.*, 136:164304, 2012; [3] Y. Sato, H. Kono *et al.*, *J. Am. Chem. Soc.*, 125:8019-8031, 2003; [4] M. Elstner, *et al.*, *Phys. Rev. B*, 58:7260-7268, 1998; [5] Experiment: J.-H. Kim *et al.*, *Phys. Rev. B*, 86:161415(R), 2012.